

## $\zeta$ -determinant and the Quillen determinant on the Grassmannian of elliptic self-adjoint boundary conditions

Simon SCOTT <sup>a</sup>, Krzysztof WOJCIECHOWSKI <sup>b</sup>

<sup>a</sup> Department of Mathematics, King's College, Strand, London WC2R 2LS, UK  
E-mail: sscott@mth.kcl.ac.uk

<sup>b</sup> Department of Mathematics, IUPUI, Indianapolis, IN 46202, USA  
E-mail: kwojciechowski@math.iupui.edu

(Reçu le 10 octobre 1998, accepté le 19 octobre 1998)

---

**Abstract.** We announce a proof of the equality of the  $\zeta$ -determinant of a Dirac operator over an odd-dimensional manifold with boundary with the Quillen determinant section computed in a canonical trivialization over the Grassmannian of elliptic self-adjoint boundary conditions. © Académie des Sciences/Elsevier, Paris

### *Le $\zeta$ -déterminant et le déterminant de Quillen sur l'espace grassmannien de conditions au bord elliptiques auto-adjointes*

**Résumé.** Dans cette Note, nous annonçons la preuve de l'égalité du  $\zeta$ -déterminant d'un opérateur de Dirac sur une variété à bord, compacte, de dimension impaire, avec la section déterminant de Quillen calculée dans une trivialisation canonique sur l'espace du grassmannien aux conditions de bords elliptiques auto-adjointes. © Académie des Sciences/Elsevier, Paris

---

### *Version française abrégée*

Soit  $\mathcal{D} : C^\infty(M; \mathcal{S}) \rightarrow C^\infty(M; \mathcal{S})$  un opérateur de Dirac compatible avec une variété  $M$  à bord, compacte, de dimension impaire, et de bord  $Y$ . On suppose que les structures sont des structures produit au voisinage de  $Y$ . On désigne par  $\text{Gr}_\infty^*(\mathcal{D})$  l'espace grassmannien lisse réduit de toutes les projections  $P$  de  $L^2(Y; \mathcal{S}|Y)$ , de manière que  $P - P(\mathcal{D})$  soit un opérateur de noyau lisse, où  $P(\mathcal{D})$  est la projection de Calderón sur l'espace des données de Cauchy définies par  $\mathcal{D}$  (voir [2]), et de plus  $-\Gamma P \Gamma = I - P$ . Dans la dernière égalité,  $\Gamma$  désigne la multiplication de Clifford par le vecteur cotangent unité et normal au bord.

---

Note présentée par Jean-Michel BISMUT.

Soit  $\gamma_0 : H^1(M; \mathcal{S}) \rightarrow H^{1/2}(Y; \mathcal{S}_Y)$  l'opérateur de trace défini par la restriction des sections de Sobolev au bord. Pour tout  $P \in \mathcal{G}r_\infty^*(\mathcal{D})$  le problème à bord elliptique :

$$\begin{cases} \mathcal{D}_P = \mathcal{D} : \text{dom}(\mathcal{D}_P) \rightarrow L^2(M; \mathcal{S}), \\ \text{dom}(\mathcal{D}_P) = \{\psi \in H^1(M; \mathcal{S}) : P\gamma_0\psi = 0\}, \end{cases}$$

définit un opérateur de Fredholm auto-adjoint de spectre réel et discret, et tel que  $\text{Ker}(\mathcal{D}_P) \subset C^\infty(M; \mathcal{S})$  (voir [2]).

Le  $\zeta$ -déterminant  $\det_\zeta \mathcal{D}_P \in \mathbf{C}$  est bien défini sur  $\mathcal{G}r_\infty^*(\mathcal{D})$  (voir [10]). Sans faire de choix, la « fonction » déterminant sur  $\mathcal{G}r_\infty^*(\mathcal{D})$  apparaît comme une section  $\det$  du fibré en doites déterminant de Quillen associé à la famille  $\{\mathcal{D}_P | P \in \mathcal{G}r_\infty^*(\mathcal{D})\}$ . Toutefois, sur  $\mathcal{G}r_\infty^*(\mathcal{D})$  ce fibré a une trivialisatation canonique qui permet de définir  $\det$  comme une fonction complexe  $\det_C : \mathcal{G}r_\infty^*(\mathcal{D}) \rightarrow \mathbf{C}$ , appelée le « déterminant canonique » (voir [7], [8]).

Le résultat principal de cette Note est l'égalité projective de ces deux régularisations du déterminant de  $\mathcal{D}_P$  :

THÉORÈME 0.1. – Sur  $\mathcal{G}r_\infty^*(\mathcal{D})$  on a

$$\det_\zeta \mathcal{D}_P = \det_\zeta \mathcal{D}_{P(\mathcal{D})} \cdot \det_C \mathcal{D}_P. \quad (1)$$

Remarque 0.2. – Dans la monographie [2], on a donné un exposé détaillé de la relation entre l'indice de l'opérateur non auto-adjoint  $\mathcal{D}_P$ , où  $P$  désigne un projecteur pseudo-différentiel quelconque de manière que  $P - P(\mathcal{D})$  soit un opérateur lissant, et l'indice de l'opérateur correspondant sur le bord  $PP(\mathcal{D}) : \text{Ran } P(\mathcal{D}) \rightarrow \text{Ran } P$ . On a montré que

$$\text{index } \mathcal{D}_P = \text{index } PP(\mathcal{D}) \quad (2)$$

(ceci a été montré pour la première fois dans [1]), c'est-à-dire que l'indice d'un opérateur agissant sur toute la variété est égal à l'indice d'un opérateur agissant seulement sur le bord. Le théorème 0.1 dit que le  $\zeta$ -déterminant de l'opérateur  $\mathcal{D}_P$  est, abstraction faite d'une constante naturelle multiplicative, égal au déterminant canonique de l'opérateur  $\mathcal{D}_P$ , qui est vraiment le déterminant de Fredholm de l'opérateur  $PP(\mathcal{D})$ . Dans ce sens, notre nouveau résultat étend (2) à une correspondance entre déterminants, même si le déterminant est un invariant beaucoup plus subtil.

Pour montrer le théorème 0.1, nous étudions la variation de  $\det_\zeta \mathcal{D}_P$  et  $\det_C \mathcal{D}_P$  par changement de la condition au bord  $P$ . Nous fixons  $P_1, P_2 \in \mathcal{G}r_\infty^*(\mathcal{D})$  et nous considérons des chemins  $P_{i,r} = U_r P_i U_r^{-1}$ , où  $\{U_r\}$  est une famille d'opérateurs elliptiques et unitaires. Nous étudions la différence

$$\log \det_\zeta \mathcal{D}_{P_{2,r}} - \log \det_\zeta \mathcal{D}_{P_{1,r}}.$$

Le résultat essentiel est :

PROPOSITION 0.3

$$\frac{d}{dr} \Big|_{r=0} \{ \log \det_\zeta \mathcal{D}_{P_{2,r}} - \log \det_\zeta \mathcal{D}_{P_{1,r}} \} = \frac{d}{dr} \Big|_{r=0} \{ \log \det_C \mathcal{D}_{P_{2,r}} - \log \det_C \mathcal{D}_{P_{1,r}} \}. \quad (3)$$

En intégrant (3), nous obtenons

$$\frac{\det_\zeta \mathcal{D}_{U P_2 U^{-1}} / \det_C \mathcal{D}_{U P_2 U^{-1}}}{\det_\zeta \mathcal{D}_{P_2} / \det_C \mathcal{D}_{P_2}} = \frac{\det_\zeta \mathcal{D}_{U P_1 U^{-1}} / \det_C \mathcal{D}_{U P_1 U^{-1}}}{\det_\zeta \mathcal{D}_{P_1} / \det_C \mathcal{D}_{P_1}} := \alpha(U).$$

Ici  $\alpha := \alpha_{P_1, P_2} : \mathcal{U}_\infty \rightarrow \mathbf{C}^*$  est un caractère du groupe  $\mathcal{U}_\infty$ . En substituant  $P_1 = P(\mathcal{D})$  et  $P_2 = P$ , nous avons  $\det_\zeta \mathcal{D}_P = \det_\zeta \mathcal{D}_{P(\mathcal{D})} \cdot \det_C \mathcal{D}_P \cdot \alpha(U)$ .

Une étude plus approfondie de  $\alpha(U)$  montre que  $\alpha = 1$ .

## 1. Introduction and statement of results

The  $\zeta$ -determinant of a self-adjoint elliptic operator has been of central importance in spectral geometry, topology, and quantum field theory (QFT). In this paper we show that the  $\zeta$ -determinant of an elliptic boundary problem on an odd-dimensional manifold is equal (up to a multiplicative constant) to the Fredholm determinant of a canonically associated operator over the boundary. This opens the door to further investigation of pasting formulae for the  $\zeta$ -determinant, as suggested by recent studies in QFT.

Let  $M$  be a compact odd-dimensional Riemannian manifold with boundary  $Y := \partial M$ , and let  $S$  denote a vector bundle of Clifford modules over  $M$ . Assume that we have a compatible Dirac operator  $\mathcal{D} : C^\infty(M; S) \rightarrow C^\infty(M; S)$  acting on sections of  $S$ . In this paper we consider only the case of product structures in a neighbourhood of the boundary  $Y$ . This means that there exists  $\mathcal{N} = [0, 1] \times Y$ , a collar neighbourhood of the boundary  $Y = \{0\} \times Y$ , such that the Riemannian metric on  $M$  and Hermitian metric on  $S$  are pull-backs of the metric structures over the boundary. Over  $\mathcal{N}$  the Dirac operator  $\mathcal{D}$  then takes the form  $\mathcal{D}|_{\mathcal{N}} = \Gamma(\partial_u + B)$ , where  $B : C^\infty(Y; S|_Y) \rightarrow C^\infty(Y; S|_Y)$  is the corresponding Dirac operator over  $Y$ , and  $\Gamma : S|_Y \rightarrow S|_Y$  is Clifford multiplication by the normal vector  $du$ . To avoid technical trivialities we assume that  $\text{Ker}(B) = \{0\}$ . Both  $B$  and  $\Gamma$  are independent of the normal coordinate  $u$  and they satisfy the identities

$$\Gamma^2 = -I, \quad \Gamma^* = -\Gamma, \quad \Gamma \cdot B + B \cdot \Gamma = 0.$$

The boundary Clifford bundle splits  $S|_Y = S^+ \oplus S^-$  into the  $\pm i$ -eigenspaces of the anti-involution  $\Gamma$ , and  $\mathcal{D}|_{\mathcal{N}}$  has the representation

$$\mathcal{D}|_{\mathcal{N}} = \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \cdot \left( \frac{\partial}{\partial u} + \begin{pmatrix} 0 & B^+ \\ B^+ & 0 \end{pmatrix} \right),$$

where  $B^\pm : F^\pm \rightarrow F^\mp$  are the chiral Dirac operators over the boundary acting on  $F^\pm = C^\infty(Y; S^\pm)$ .

Let  $\Pi_{>}$  be the spectral projection on  $L^2(Y; S|_Y)$  with range equal to the subspace  $H^+$  spanned by the eigensections of  $B$  with positive eigenvalues. We consider elliptic self-adjoint boundary conditions for  $\mathcal{D}$  parameterized by the smooth restricted Grassmannian  $\text{Gr}_\infty^*(\mathcal{D})$  of all orthogonal projections  $P$  on  $L^2(Y; S|_Y)$  such that  $P - \Pi_{>}$  is an operator with a smooth kernel and such that  $-\Gamma P \Gamma = I - P$ . For each  $P \in \text{Gr}_\infty^*(\mathcal{D})$  the associated elliptic boundary value problem

$$\mathcal{D}_P := \mathcal{D} : \text{dom}(\mathcal{D}_P) \rightarrow L^2(M; S)$$

with  $\text{dom}(\mathcal{D}_P) = \{\psi \in H^1(M; S) : P(\psi|_Y) = 0\}$  is a self-adjoint Fredholm operator with  $\text{Ker}(\mathcal{D}_P) \subset C^\infty(M; S)$  and real discrete spectrum (cf. [2]).

The space  $\text{Gr}_\infty^*(\mathcal{D})$  is characterized by two important properties. First, it contains the Calderón projection  $P(\mathcal{D})$  with range equal to  $\mathcal{H}(\mathcal{D})$ , the space of Cauchy data of the operator  $\mathcal{D}$ :

$$\mathcal{H}(\mathcal{D}) = \{f \in C^\infty(Y; S|_Y) : \exists s \in C^\infty(M; S), \mathcal{D}s = 0, s|_Y = f\}.$$

Second,  $H^+ = \text{range}(\Pi_{>})$  is the graph of the elliptic unitary isomorphism  $(B^+ B^-)^{-1/2} B^+ : F^+ \rightarrow F^-$ , the “phase” of  $B^+$ . Similarly, any other elliptic boundary condition  $P \in \text{Gr}_\infty^*(\mathcal{D})$  is described precisely by the property that it is the graph of an elliptic unitary isomorphism  $T : F^+ \rightarrow F^-$  such that  $T - (B^+ B^-)^{-1/2} B^+$  has a smooth kernel. In particular, there is a specific unitary isomorphism  $K$  such that  $\mathcal{H}(\mathcal{D}) = \text{graph}(K : F^+ \rightarrow F^-)$ . The graph characterization induces a canonical trivialization  $\sigma_{\text{graph}}$  of the determinant line bundle  $\mathcal{L}$  over  $\text{Gr}_\infty^*(\mathcal{D})$  associated to the family of elliptic boundary value problems  $\{\mathcal{D}_P : P \in \text{Gr}_\infty^*(\mathcal{D})\}$  (see [7], [8]). Computing the ratio of the determinant section  $P \rightarrow \det \mathcal{D}_P$  of  $\mathcal{L}$  (the Quillen determinant) to  $\sigma_{\text{graph}}$  provides us with a canonical regularization of the determinant of  $\mathcal{D}_P$  as a complex number  $\det_c \mathcal{D}_P$  called the canonical determinant. One

finds (see [7], [8])

$$\det_C \mathcal{D}_P = \det_{F^+} \left[ \frac{1}{2}(I + KT^{-1}) \right], \quad (1)$$

where  $\text{range}(P) = \text{graph}(T : F^+ \rightarrow F^-)$ . Here the right-hand side of (1) is the Fredholm determinant taken on  $F^-$  and since  $KT^{-1}$  has the form  $I + \text{trace-class}$  it is well defined.

On the other hand, for  $P \in \mathcal{G}r_\infty^*(\mathcal{D})$  the self-adjoint elliptic operator  $\mathcal{D}_P$  has a well-defined  $\zeta$ -function regularized determinant

$$\det_\zeta \mathcal{D}_P := e^{i\frac{\pi}{2}(\eta_{\mathcal{D}_P}(0) - \zeta_{\mathcal{D}_P^2}(0))} \cdot e^{i\frac{1}{2}\zeta'_{\mathcal{D}_P^2}(0)}.$$

The existence of the meromorphic extensions of the functions  $\zeta_{\mathcal{D}_P^2}(s)$  and  $\eta_{\mathcal{D}_P}(s)$  to the whole complex plane, and their nice behaviour in a neighbourhood of  $s = 0$  was established in the paper [10]. While the canonical determinant is a robust algebraic regularization which is relatively easy to compute, the  $\zeta$ -determinant is, in contrast, a highly delicate analytic regularization which is notoriously difficult to compute and there is little a priori reason to suppose that the two regularizations might coincide. The main result of this announcement is the following remarkable fact:

**THEOREM 1.1.** – *The following equality holds over  $\mathcal{G}r_\infty^*(\mathcal{D})$*

$$\det_\zeta \mathcal{D}_P = \det_\zeta \mathcal{D}_{P(\mathcal{D})} \cdot \det_C \mathcal{D}_P. \quad (2)$$

That is,

$$\det_\zeta \mathcal{D}_P = c \cdot \det_{F^+} \left[ \frac{1}{2}(I + KT^{-1}) \right], \quad (3)$$

where the constant  $c$  is the value of the  $\zeta$ -determinant evaluated at the Cauchy data space (Calderón) boundary condition.

Equivalently, (2) expresses the equality of the ratios  $\det_\zeta \mathcal{D}_P / \det_\zeta \mathcal{D}_{P(\mathcal{D})}$  and  $\det_C \mathcal{D}_P / \det_C \mathcal{D}_{P(\mathcal{D})}$ .

**Remark 1.2.** – In the monograph [2] a detailed account was given of the relation between the index of the operator  $\mathcal{D}_P$  for any  $P$  in the smooth Grassmannian  $\mathcal{G}r_\infty(\mathcal{D})$  of projections differing from  $\Pi_{>}$  by a smoothing operator, and the index of the boundary operator  $PP(\mathcal{D}) : \mathcal{H}(\mathcal{D}) \rightarrow \text{Ran } P$ . Outside of the real submanifold  $\mathcal{G}r_\infty^*(\mathcal{D})$  of  $\mathcal{G}r_\infty(\mathcal{D})$  the Fredholm operator  $\mathcal{D}_P$  is not self-adjoint and it was shown that  $\text{index } \mathcal{D}_P = \text{index } PP(\mathcal{D})$  (see [1] for the original presentation). Hence  $\text{index } \mathcal{D}_P$  is equal to the index of an operator which lives only on the boundary. Theorem 1.1 extends this somewhat unexpectedly to a correspondence between the determinants.

The proof of Theorem 1.1 is outlined in Section 2. Details will appear in [9]. Note that (3) generalizes well-known formulas for the  $\zeta$ -determinant in dimension one (see in particular [3]).

## 2. Sketch of the proof

The proof of Theorem 1.1 adapts ideas suggested in [5] by Forman who obtained precise results in one and two dimensions. We show that this method generalizes to compute  $\det_\zeta \mathcal{D}_P$  over  $\mathcal{G}r_\infty^*(\mathcal{D})$  in any odd dimension.

To establish Theorem 1.1 we study the variation of  $\det_\zeta \mathcal{D}_P$  and  $\det_C \mathcal{D}_P$  under a change of the boundary condition  $P$ . To compute the variation we exploit the transitive action  $P \mapsto UPU^{-1}$  on  $\mathcal{G}r_\infty^*(\mathcal{D})$  of the group  $\mathcal{U}_\infty$  of unitary isomorphisms of the form:

$$U = \begin{pmatrix} I_{F^+} & 0 \\ 0 & g \end{pmatrix}, \quad (4)$$

written relative to the grading  $F^+ \oplus F^-$ , where  $g - I_{F^-}$  is a smoothing operator, with  $I_{F^\pm}$  the identity operator on  $F^\pm$ . Using a “unitary twist” this action is extended to a unitary action on  $C^\infty(M; \mathcal{S})$  which by spectral invariance allows us to compute the variation of  $\det_\zeta \mathcal{D}_P$  with respect to  $P$  (see [4], [8], [10] for details). Fix  $P_1, P_2 \in \mathcal{G}r_\infty^*(\mathcal{D})$  and consider the paths  $P_{i,r} = U_r P_i U_r^{-1}$  determined by a smooth 1-parameter family  $U_r \in \mathcal{U}_\infty$  with  $-1 < r < 1$  and  $U_0 = I$ . We study the difference

$$\log \det_\zeta \mathcal{D}_{P_{2,r}} - \log \det_\zeta \mathcal{D}_{P_{1,r}}.$$

The crucial result is:

PROPOSITION 2.1. – *The following equality holds:*

$$\frac{d}{dr} \Big|_{r=0} \{ \log \det_\zeta \mathcal{D}_{P_{2,r}} - \log \det_\zeta \mathcal{D}_{P_{1,r}} \} = \frac{d}{dr} \Big|_{r=0} \{ \log \det_C \mathcal{D}_{P_{2,r}} - \log \det_C \mathcal{D}_{P_{1,r}} \}. \quad (5)$$

*Outline proof of Proposition 2.1.* – We start with the discussion of the left side of equation (5). We need the following result proved in [10].

PROPOSITION 2.2. – (1). *The value of the  $\zeta$ -function of  $\mathcal{D}_P$  at  $s = 0$  is constant on  $\mathcal{G}r_\infty^*(\mathcal{D})$ .*

(2). *The variation of the  $\eta$ -invariant of  $\mathcal{D}_P$  is stable in the following sense: for any pair of projections  $P_1, P_2 \in \mathcal{G}r_\infty^*(\mathcal{D})$*

$$\frac{d}{dr} \Big|_{r=0} (\eta_{\mathcal{D}_{P_{1,r}}}(0)) = \frac{d}{dr} \Big|_{r=0} (\eta_{\mathcal{D}_{P_{2,r}}}(0)),$$

where the families  $\{P_{i,r} = U_r P_i U_r^{-1}\}_{r \in (-1,1)}$  are constructed by using a family of unitary operators  $\{U_r\}$  of the form (4), such that  $g_r$  is a smooth family of unitary operators on  $F^-$  of the form  $I +$  smoothing operator with  $g_0 = I$ .

It follows from Proposition 2.2 that we have only to study the variation of the modulus  $\det_\zeta [\mathcal{D}_{P_{i,r}}] = e^{-\frac{1}{2} \zeta'_{\mathcal{D}_{i,r}}(0)}$  of the determinant  $\det_\zeta \mathcal{D}_{P_{i,r}}$ . We compute:

$$\frac{d}{dr} \Big|_{r=0} \left( \frac{1}{2} \zeta'_{\mathcal{D}_{i,r}}(0) \right) = - \lim_{\varepsilon \rightarrow 0} \text{Tr} \left( \frac{d}{dr} \Big|_{r=0} \mathcal{D}_{P_{i,r}} \right) \mathcal{D}_{P_{i,r}}^{-1} e^{-\varepsilon \mathcal{D}_{P_{i,r}}^2}. \quad (6)$$

To evaluate the limit in (7) we have to study the inverse of the operator  $\mathcal{D}_P$ .

The operator  $\mathcal{D}_P$  is invertible if and only if the operator  $PP(\mathcal{D}) : \mathcal{H}(\mathcal{D}) \rightarrow \text{range}(P)$  is invertible, and in this case we have:

PROPOSITION 2.3. – *The operator inverse  $\mathcal{D}_P^{-1} : L^2(M; \mathcal{S}) \rightarrow \text{dom}(\mathcal{D}_P)$  of the operator  $\mathcal{D}_P$  is given by the formula:*

$$\mathcal{D}_P^{-1} = \mathcal{D}^{-1} - \mathcal{D}^{-1} \gamma_0^* \Gamma(PP(\mathcal{D}))^{-1} P \gamma_0 \mathcal{D}^{-1}, \quad (7)$$

where  $\mathcal{D}^{-1}$  denotes the operator  $r_M \cdot \tilde{\mathcal{D}}^{-1} \cdot e_{\tilde{M}}$ . Here  $\tilde{\mathcal{D}}$  is the (invertible) double operator of  $\mathcal{D}$  over the closed double  $\tilde{M}$  of  $M$ ;  $e_{\tilde{M}} : H^s(M; \mathcal{S}) \rightarrow H^s(\tilde{M}; \tilde{\mathcal{S}})$  and  $r_M : H^s(\tilde{M}; \tilde{\mathcal{S}}) \rightarrow H^s(M; \mathcal{S})$  are the corresponding extension and restriction operators; and  $\gamma_0^*$  is the dual to the operator  $\gamma_0 : H^s(M; \mathcal{S}) \rightarrow H^{s-1/2}(Y; \mathcal{S}|Y)$  restricting sections to the boundary (see [2] for details).

Note that  $(PP(\mathcal{D}))^{-1}$  is the restriction to  $\text{range}(P)$  of the elliptic pseudodifferential operator  $((I - P)(I - P(\mathcal{D})) + PP(\mathcal{D}))^{-1}$ . We observe that the difference

$$\mathcal{D}_{P_1}^{-1} - \mathcal{D}_{P_2}^{-1} = \mathcal{D}^{-1} \gamma^* \Gamma(P_1 P(\mathcal{D}))^{-1} P_1 \gamma_0 \mathcal{D}^{-1} - \mathcal{D}^{-1} \gamma^* \Gamma(P_2 P(\mathcal{D}))^{-1} P_2 \gamma_0 \mathcal{D}^{-1},$$

is a smoothing operator, which implies the next Proposition.

S. Scott, K. Wojciechowski

PROPOSITION 2.4. – For any pair of projections  $P_1, P_2 \in \mathcal{G}r_\infty^*(\mathcal{D})$  such that  $\mathcal{D}_{P_1}$  and  $\mathcal{D}_{P_2}$  are invertible operators and any 1-parameter family of unitary operators  $U_r$  of the form (4) the following identity holds:

$$\frac{d}{dr} \Big|_{r=0} \{ \log \det_\zeta \mathcal{D}_{P_{2,r}} - \log \det_\zeta \mathcal{D}_{P_{1,r}} \} = \text{Tr} \dot{D}(D_{P_2}^{-1} - D_{P_1}^{-1}), \quad (8)$$

where  $\dot{D}$  denotes the operator  $\frac{d}{dr}(U_r^{-1} \mathcal{D} U_r)|_{r=0}$ , and the families  $P_{i,r}$  are defined as in Proposition 2.2.

The determinant lines of the operators  $PP(\mathcal{D}) : \mathcal{H}(\mathcal{D}) \rightarrow \text{Ran}(P)$  and  $\mathcal{D}_P$  are canonically isomorphic and under the isomorphism their abstract Quillen determinant elements identified. Hence,  $\det_C(PP(\mathcal{D}))$  is also defined by the right-side of (1), which gives us

$$\det_C \mathcal{D}_{P_{2,r}} (\det_C \mathcal{D}_{P_{1,r}})^{-1} = \det_C(P_{2,r}P(\mathcal{D})) (\det_C(P_{1,r}P(\mathcal{D})))^{-1}. \quad (9)$$

Equation (9) and the fact that  $P(\mathcal{D}) = \gamma_0 \mathcal{D}^{-1} \gamma^* \Gamma$  allow us to study the right-side of (5). Using formula (7) for the operator  $\mathcal{D}_P^{-1}$  we compute the logarithm variation of the right-side of (9) as:

PROPOSITION 2.5.

$$\frac{d}{dr} \Big|_{r=0} \log \det_C(P_{2,r} \gamma_0 \mathcal{D}^{-1} \gamma^* \Gamma(P_{1,r} P(\mathcal{D}))^{-1} P_{1,r}) = \text{Tr} \dot{\mathcal{D}}(\mathcal{D}_{P_{2,r}}^- - \mathcal{D}_{P_{1,r}}^-). \quad (10)$$

This completes the proof of Proposition 2.1.

The subspace of  $\mathcal{G}r_\infty^*(\mathcal{D})$  consisting of projections  $P$  such that  $PP(\mathcal{D}) : \mathcal{H}(\mathcal{D}) \rightarrow \text{Ran } P$  is an invertible operator is path connected (see [6], Proposition 3.12). This implies that two invertible operators  $\mathcal{D}_{P_1}$  and  $\mathcal{D}_{P_2}$  can always be connected through a path of invertible operators  $\{\mathcal{D}_{P_i}\}$ . Therefore we can integrate (5) and obtain:

$$\frac{\det_\zeta \mathcal{D}_{UP_2U^{-1}} / \det_C \mathcal{D}_{UP_2U^{-1}}}{\det_\zeta \mathcal{D}_{P_2} / \det_C \mathcal{D}_{P_2}} = \frac{\det_\zeta \mathcal{D}_{UP_1U^{-1}} / \det_C \mathcal{D}_{UP_1U^{-1}}}{\det_\zeta \mathcal{D}_{P_1} / \det_C \mathcal{D}_{P_1}} := \alpha(U).$$

For  $U, V \in \mathcal{U}_\infty$  we have  $\alpha(UV) = \alpha(U)\alpha(V)$  and hence  $\alpha : \mathcal{U}_\infty \rightarrow \mathbf{C}^*$  is a character on the group  $\mathcal{U}_\infty$ . Substituting  $P_1 = P(\mathcal{D})$ ,  $P_2 = P = UP(\mathcal{D})U^{-1}$  we obtain  $\det_\zeta \mathcal{D}_P = \det_\zeta \mathcal{D}_{P(\mathcal{D})} \cdot \det_C \mathcal{D}_P \cdot \alpha(U)$ . Further inspection of  $\alpha(U)$  as the metric in the collar neighbourhood  $\mathcal{N}$  is blownup in the normal direction to the boundary shows that  $\alpha = 1$ . This completes the proof of Theorem 1.1.

## References

- [1] Booß-Bavnbek B., Wojciechowski K.P., Desuspension of splitting elliptic symbols, *III*, Ann. Global Anal. Geom. 3/4 (1985/1986) 337–383/349–400.
- [2] Booß-Bavnbek B., Wojciechowski K.P., Elliptic Boundary Problems for Dirac Operators, Birkhäuser, Boston, 1993.
- [3] Booß-Bavnbek B., Scott S.G., Wojciechowski K.P., The  $\zeta$ -determinant and  $\mathcal{C}$ -determinant on the Grassmannian in dimension one, Lett. Math. Phys. 45 (1998) 353–362.
- [4] Douglas R.G., Wojciechowski K.P., Adiabatic limits of the  $\eta$ -invariants. The odd-dimensional Atiyah–Patodi–Singer problem, Commun. Math. Phys. 142 (1991) 139–168.
- [5] Forman R., Functional determinants and geometry, Invent. Math. 88 (1987) 447–493.
- [6] Nicolaescu L., Generalized Symplectic Geometries and the Index of Families of Elliptic Problems, Mem. Amer. Math. Soc. 609, Providence, 1997.
- [7] Scott S.G., Determinants of Dirac boundary value problems over odd-dimensional manifolds, Commun. Math. Phys. 173 (1995) 43–76.
- [8] Scott S.G., Wojciechowski K.P., Abstract Determinant and  $\zeta$ -Determinant on the Grassmannian, Lett. Math. Phys. 40 (1997) 135–145.
- [9] Scott S.G., Wojciechowski K.P., The  $\zeta$ -determinant and Quillen’s determinant for the Dirac operator on a manifold with boundary. (in preparation).
- [10] Wojciechowski K.P., The  $\zeta$ -determinant and the additivity of the  $\eta$ -invariant on the smooth, self-adjoint Grassmannian, Commun. Math. Phys. (1997) (to appear).