

Integrability - Assessment, due 17 December 2012

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Problem 1

a) Derive the fundamental Sklyanin relation

$$\{L_1(\lambda), L_2(\mu)\} = [r_{12}(\lambda, \mu), L_1(\lambda)L_2(\mu)]$$

from the local Poisson brackets

$$\{U_1(\lambda; x), U_2(\mu; y)\} = [r_{12}(\lambda, \mu), U_1(\lambda; x) + U_2(\mu; y)]\delta(x - y)$$

for the connection $U(\lambda; x)$ of the general zero-curvature formulation of integrable field theory. Recall that

$$L(\lambda) = \mathcal{P} \exp \left[\int_{-\infty}^{\infty} ds U(\lambda; s) \right].$$

b) Show that the fundamental Sklyanin relation implies that $\text{Tr}(L(\lambda))$ is a generating function for Poisson commuting quantities (you do not need to discuss locality!).

Problem 2

a) Calculate the conserved quantity

$$Q_2 = \left(\frac{d^2 F(\lambda)}{d\lambda} F(\lambda)^{-1} \right)_{\lambda=i/2} - Q_1^2$$

where

$$F(\lambda) = \text{Tr}_a T_a(\lambda)$$

and $T_a(\lambda)$ is the transfer matrix for the Heisenberg chain. You only need to bring Q_2 into a form that makes it obvious that it is a local conserved quantity. In addition to the relations from the lectures / notes, you can use the following relations:

- The conserved quantity Q_1 is

$$Q_1 = -i \sum_j P_{j,j+1}$$

where $P_{j,k}$ is the permutation operator between the sites j and k .

- The relations

$$\text{Tr}_a (P_{N,a} \cdots P_{1,a}) = U, \quad \text{Tr}_a \left(P_{N,a} \cdots \widehat{P_{j,a}} \cdots \widehat{P_{k,a}} \cdots P_{1,a} \right) = U_{j,k}$$

where U is the “clockwise” cyclic permutation operator, and $U_{j,k}$ is the clockwise cyclic permutation operator which skips the sites j and k .

- The relation

$$U_{j,k} = : P_{j,j+1} P_{k,k+1} : U$$

where the “normal ordering” of permutation operators $: P_{j,j+1} P_{k,k+1} :$ orders the factors in such a way that the second index of the first factor is not equal to the first index of the second factor.

- b) Calculate the exchange relations for the pairs $B(\lambda), A(\mu)$; $B(\lambda), B(\mu)$ and $B(\lambda), D(\mu)$ of auxiliary-space matrix elements of the transfer matrix, from the RTT relations with the XXZ-chain R matrix given by

$$R(\lambda, \mu) = \begin{pmatrix} \varphi(\lambda - \mu + \eta) & 0 & 0 & 0 \\ 0 & \varphi(\lambda - \mu) & \varphi(\eta) & 0 \\ 0 & \varphi(\eta) & \varphi(\lambda - \mu) & 0 \\ 0 & 0 & 0 & \varphi(\lambda - \mu + \eta) \end{pmatrix}$$

where

$$\varphi(\lambda) = \sinh(\lambda)$$

Problem 3

- a) Write down the Yang-Baxter equations for the two-particle scattering matrix $S_{a_1, a_2}^{b_1, b_2}(\theta_1 - \theta_2)$ of factorized scattering theory. Explain in your words, with drawings, how these Yang-Baxter equations are derived from the presence of higher-spin local conserved quantities Q_s .

- b) Consider a QFT with only one particle type. By using the resolution of the identity into asymptotic states and by using space-time translation and relativistic invariances, show that the two-point function of a hermitian, spinless field \mathcal{O} can be written, with $x^2 > t^2$, as

$$\langle \text{vac} | \mathcal{O}(x, t) \mathcal{O}(0, 0) | \text{vac} \rangle = \sum_{n=0}^{\infty} \int_{\theta_1 > \dots > \theta_n} d\theta_1 \dots d\theta_n |f(\theta_1, \dots, \theta_n)|^2 e^{-\sqrt{x^2 - t^2} \sum_j m \cosh \theta_j}.$$

Here

$$f(\theta_1, \dots, \theta_n) := \langle \text{vac} | \mathcal{O}(0, 0) | \theta_1, \dots, \theta_n \rangle$$

are the form factors of the field positioned at the space-time origin. **Hint:** use relativistic invariance of the form factors to make a shift of the rapidities by $i\pi/2$, and recall that $\sinh(\theta + i\pi/2) = i \cosh \theta$ (or, if you do the problem differently from what I have in mind, make a shift by another appropriate pure imaginary value).

References

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